

# Restoring High-Resolution Electron Spectra from the CASSIOPE/e-POP Suprathermal Electron Imager

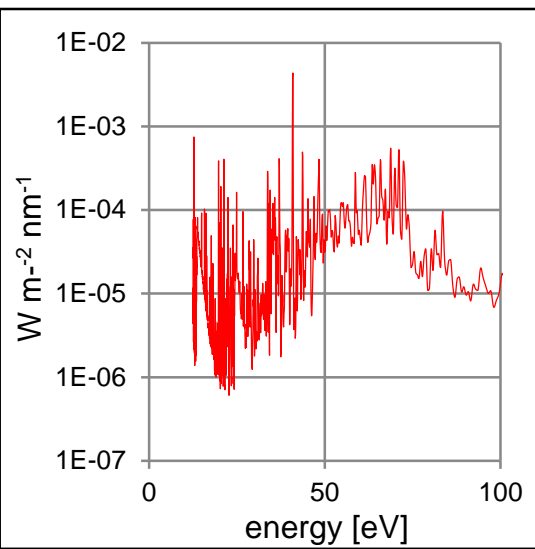
*Alexei Kouznetsov, Johnathan Burchill, David  
Knudsen*  
Department of Physics and Astronomy

## Acknowledgments:

- The work was supported by the Canadian Space Agency
- e-POP team produced a new quaternions data product to support the project

# Introduction

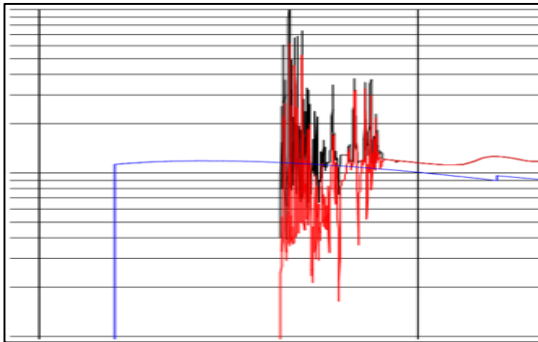
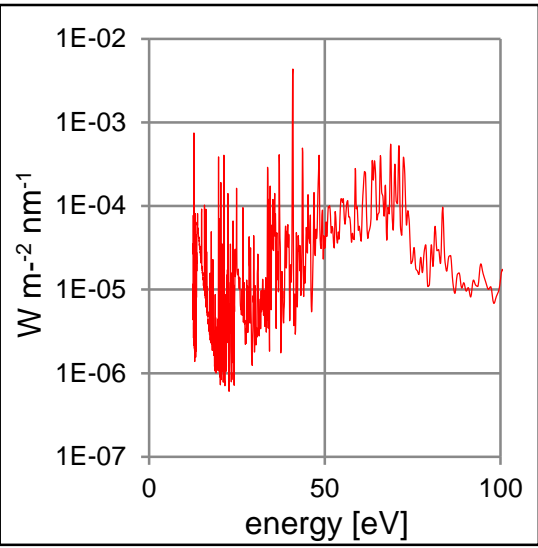
15.4 (3.8 Lyman Alpha) mW/m<sup>2</sup>



# Introduction

2

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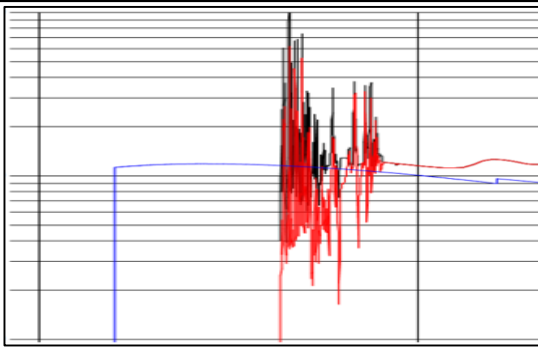
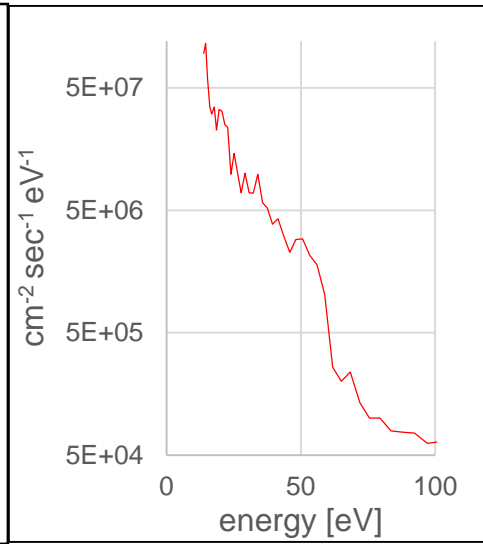
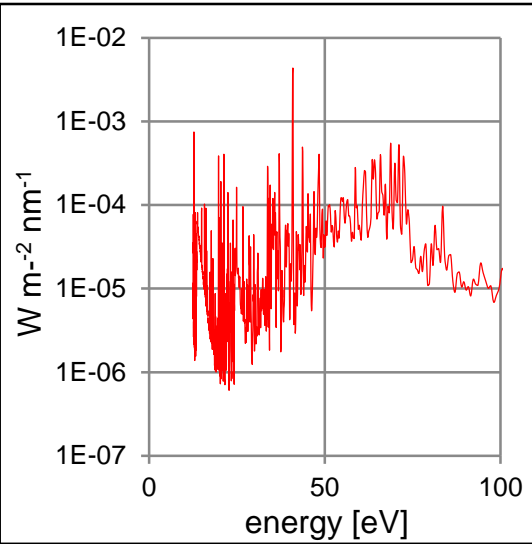


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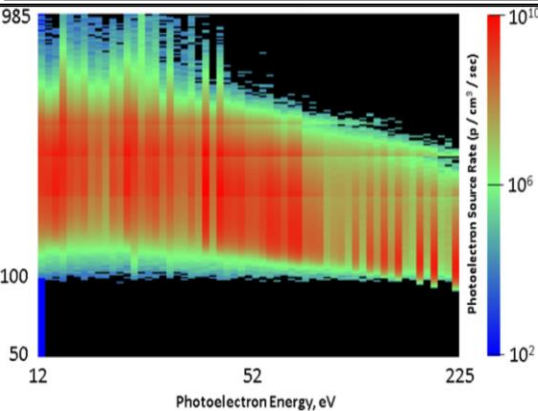
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Photoelectron Spectrum



Plasma Factory: Photoelectron Production & its Spectra Altitude Profile.

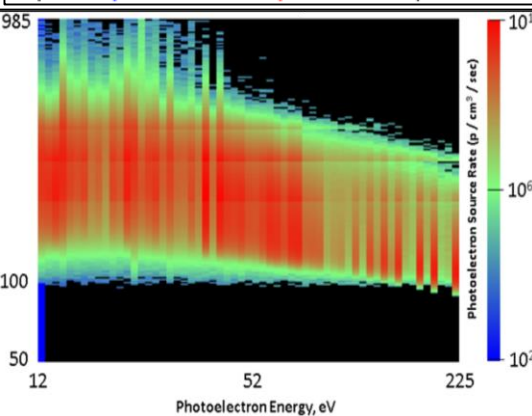
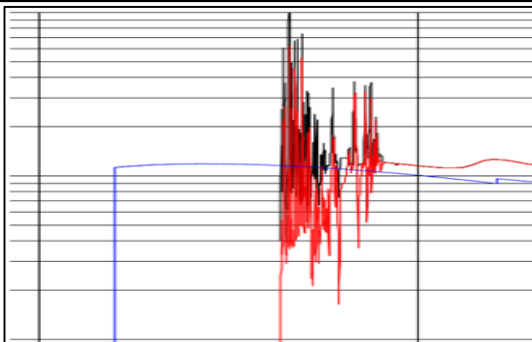
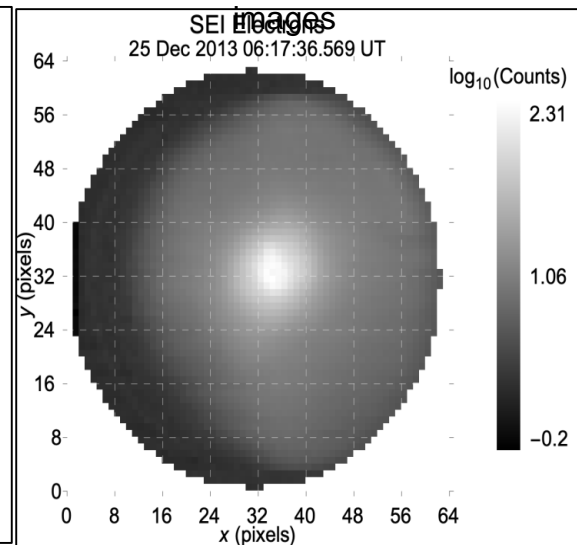
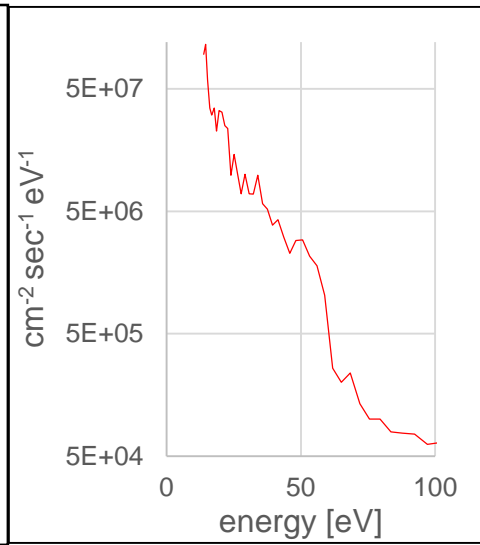
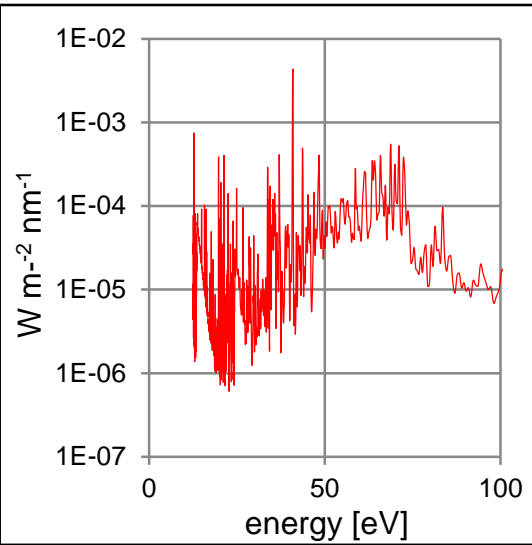


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An average of 100 SEI



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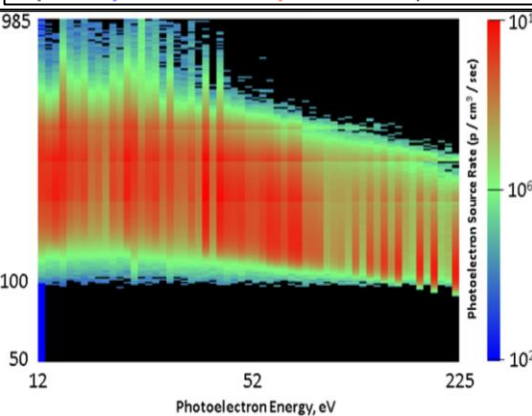
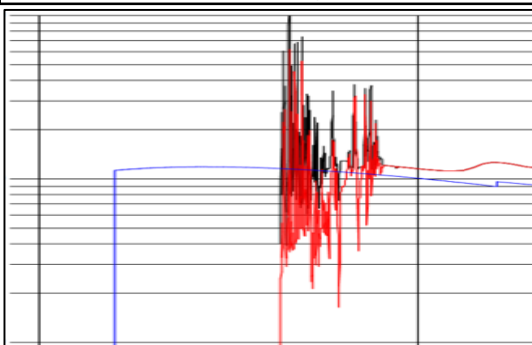
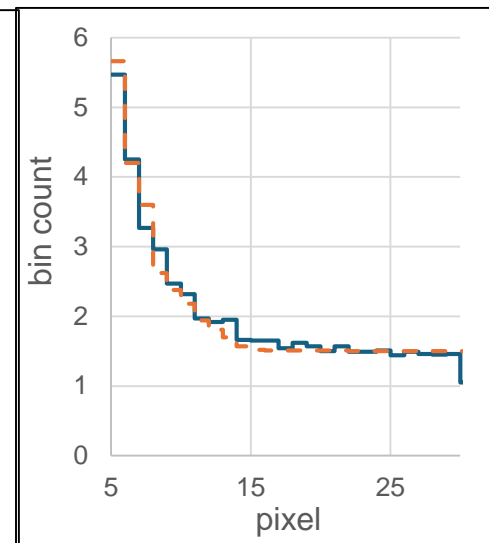
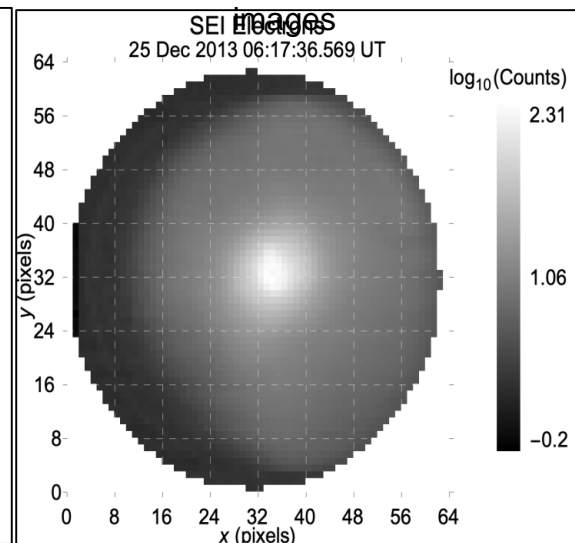
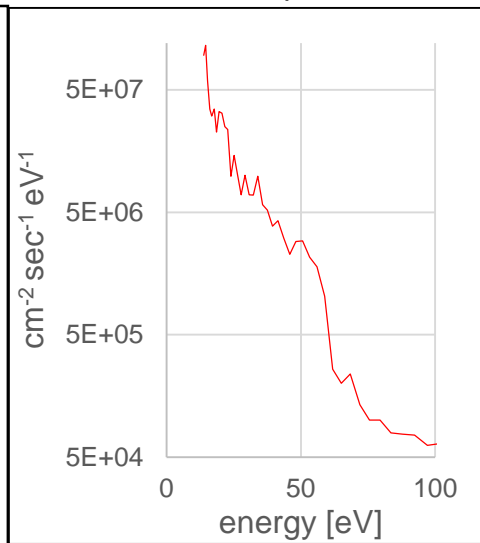
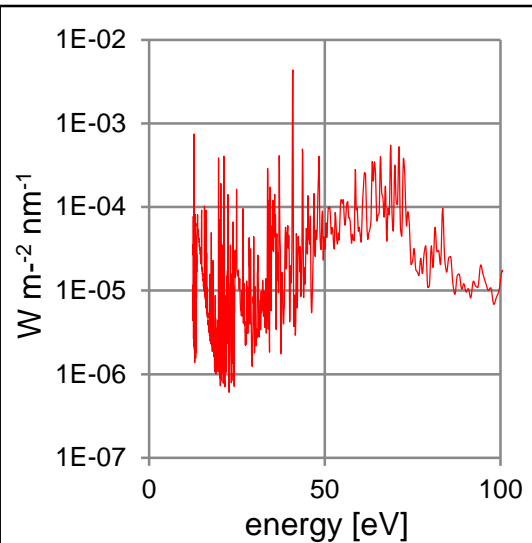
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Convolved Electron Spectra



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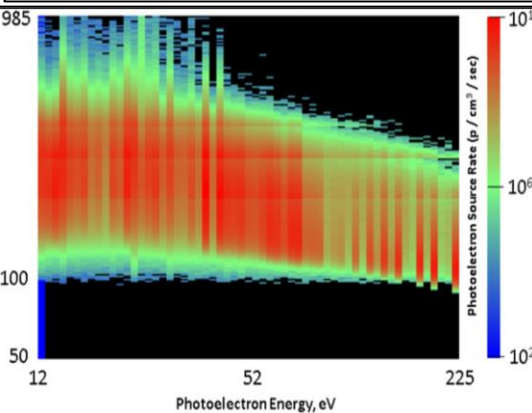
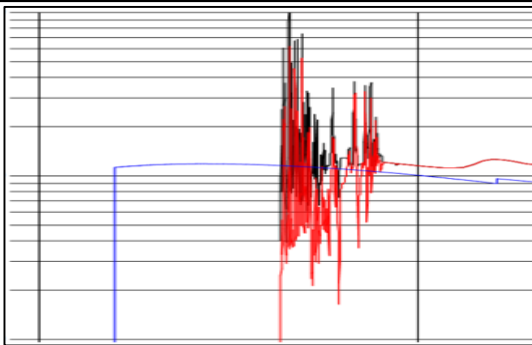
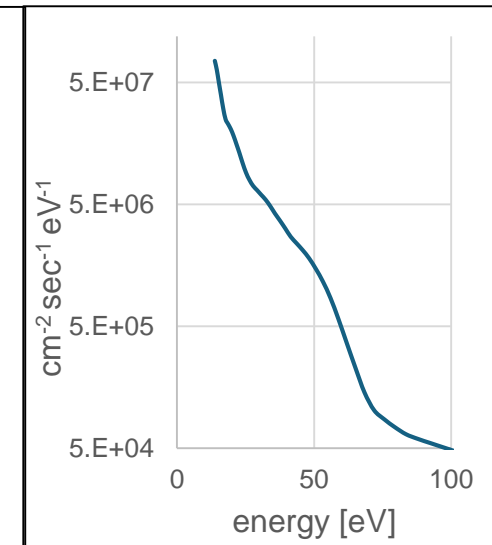
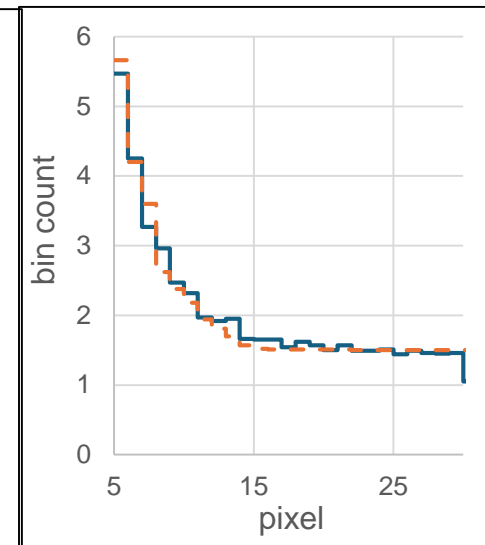
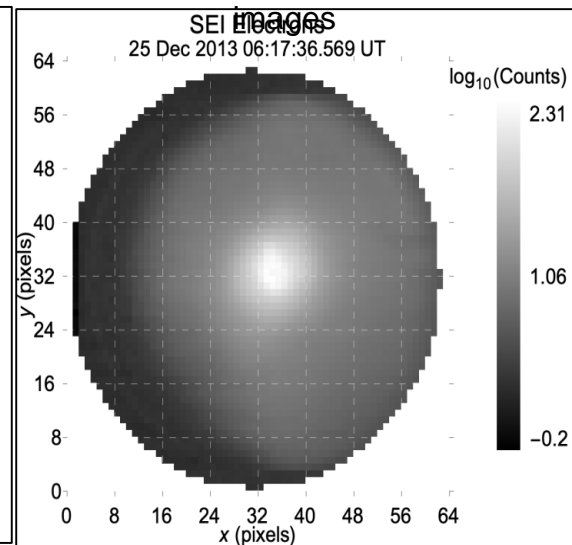
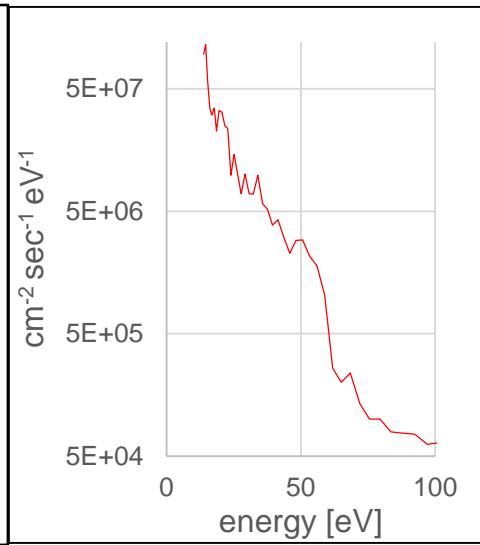
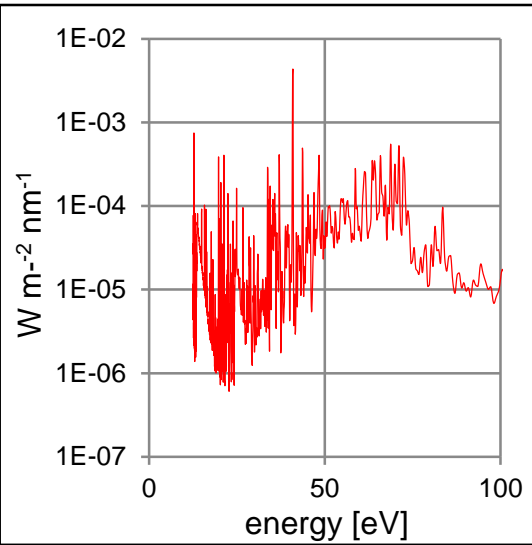
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Deconvoluted Electron Spectrum



Plasma Factory: Photoelectron Production & its Spectra Altitude Profile.

Deconvolution an Electron Spectrum from an SEI Image Based on the **Point-Spread Functions**:

1. Debye Shielding of the SEI Instrument by its high negative voltage (strongly-coupled plasma).
2. Photoelectron Collimating.
3. Photoelectron propagation in the SEI strong electric field.



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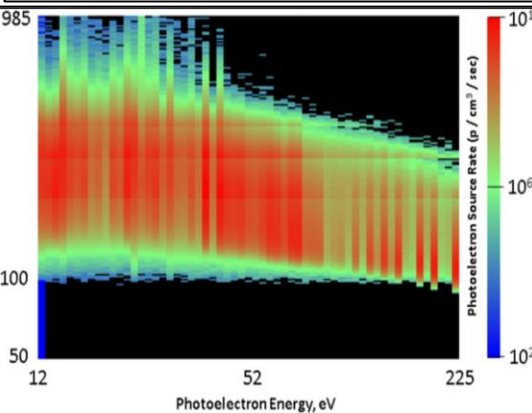
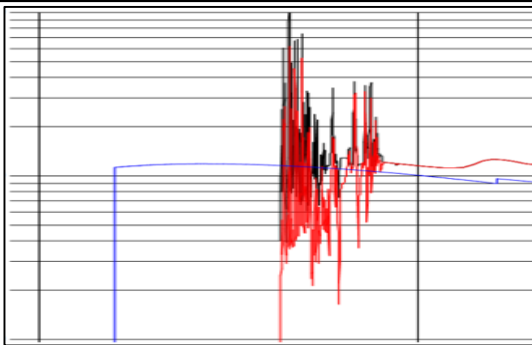
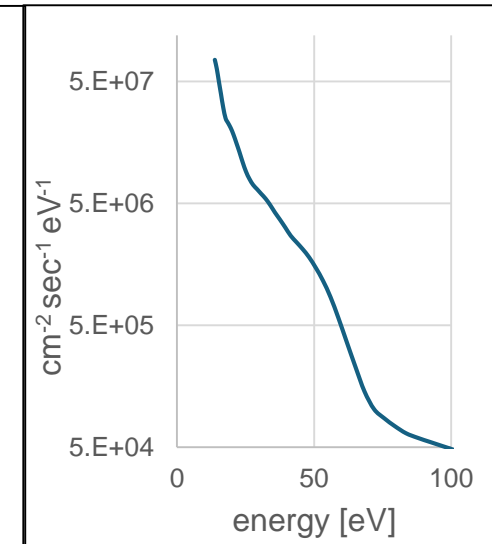
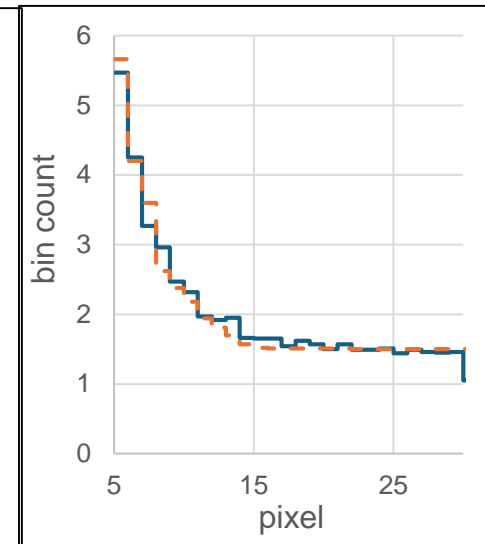
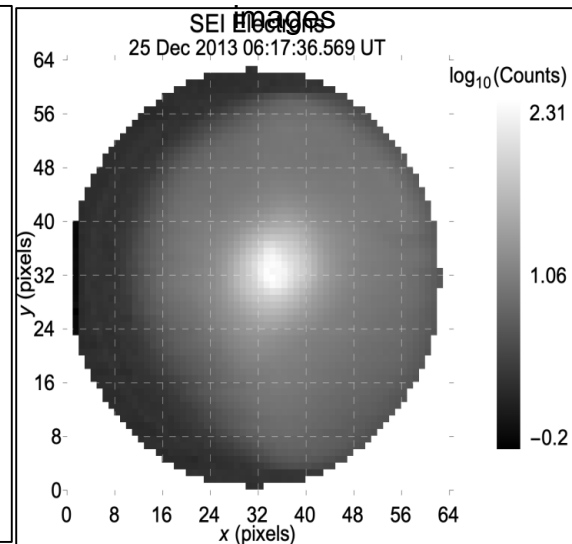
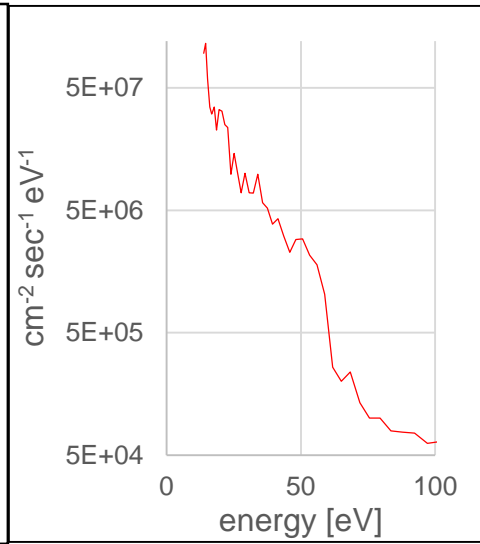
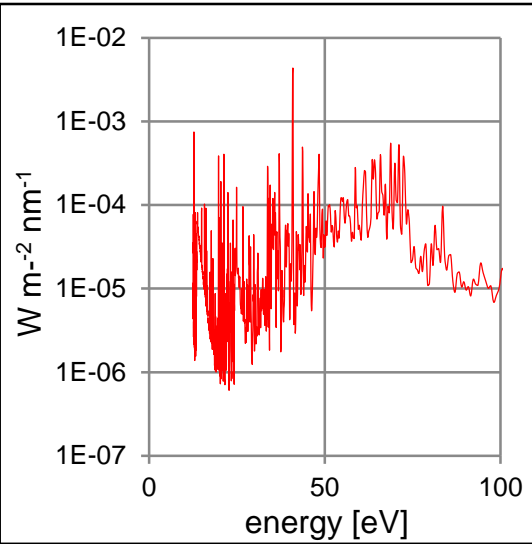
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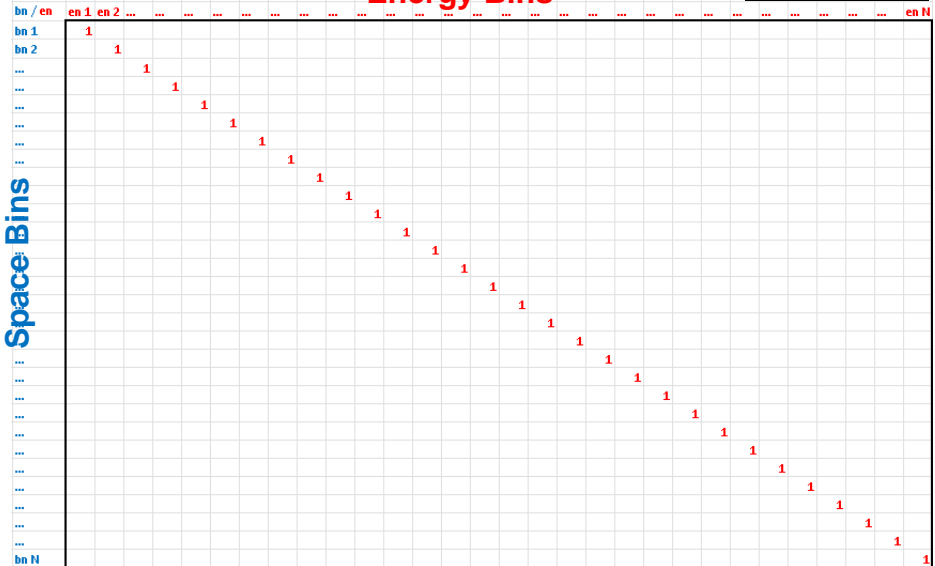
**Satellite potential is CRUCIAL**  
for the **Point-Spread Function** reliable calculations:



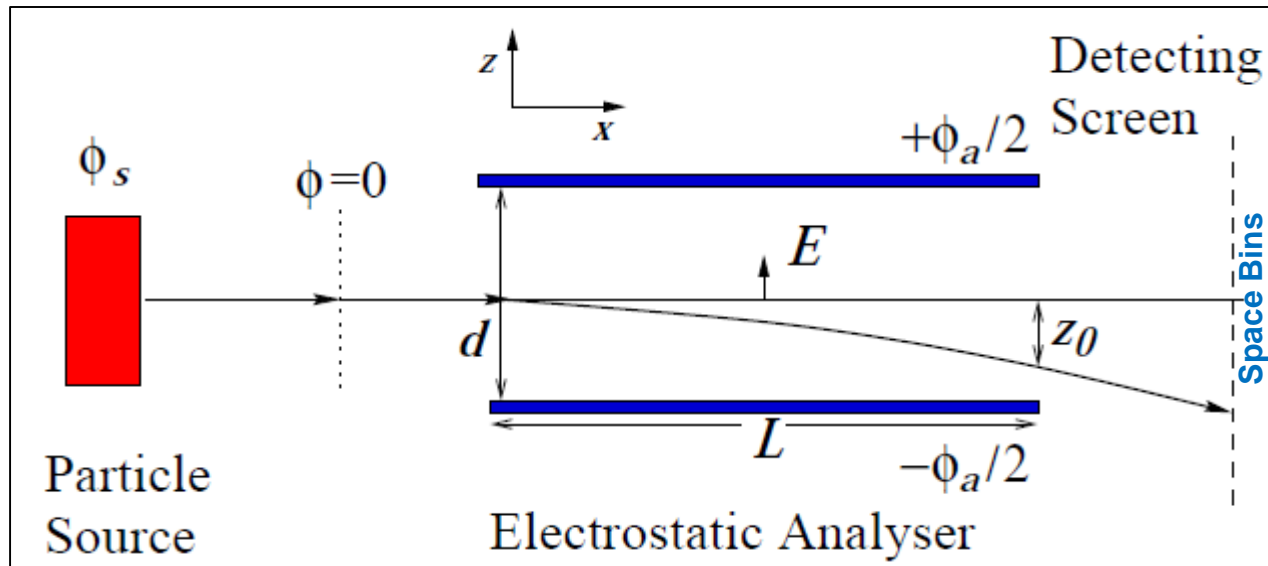
# Electrostatic Analyzer: Point-Spread Function

Energy Bins

Function



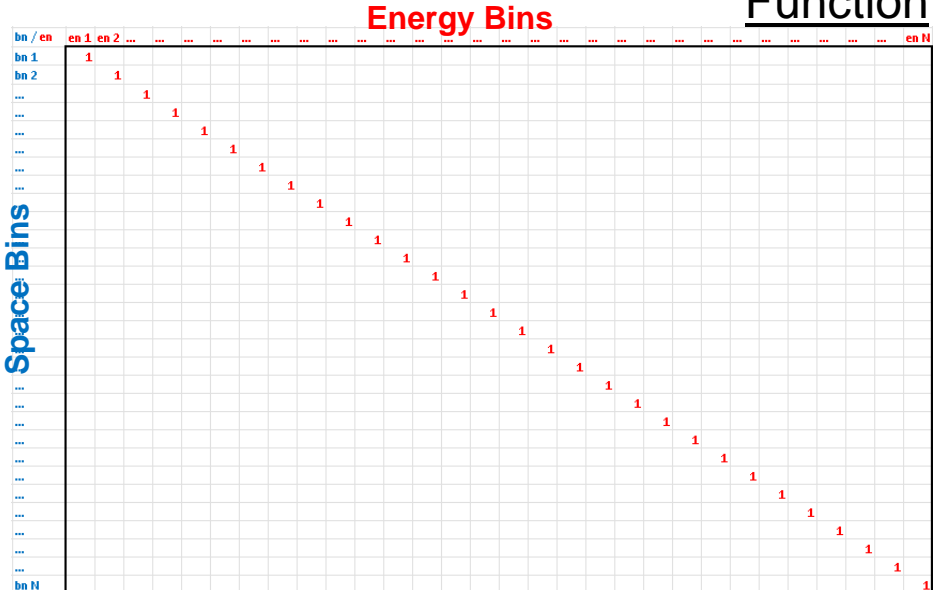
## Particle Analyzers



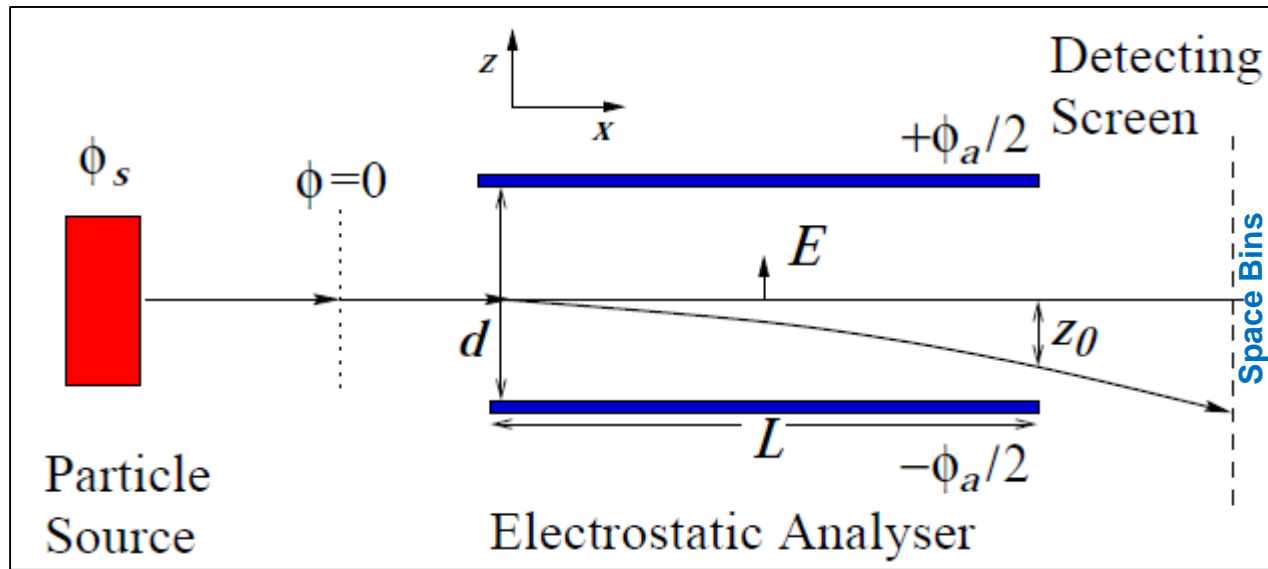
Particle Energy

$$z_0(W) = \frac{1}{4} \frac{qE}{W} L^2$$

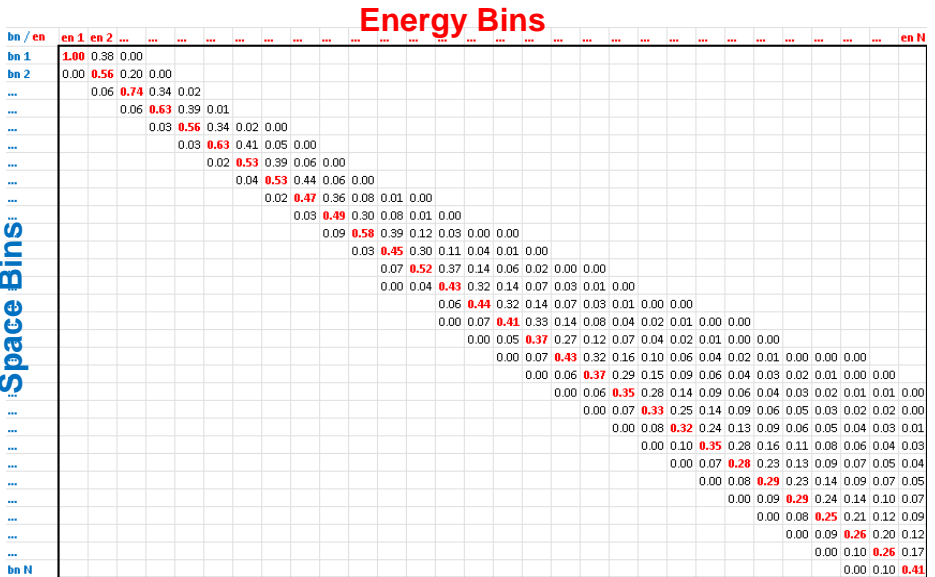
# Electrostatic Analyzer: Point-Spread Function



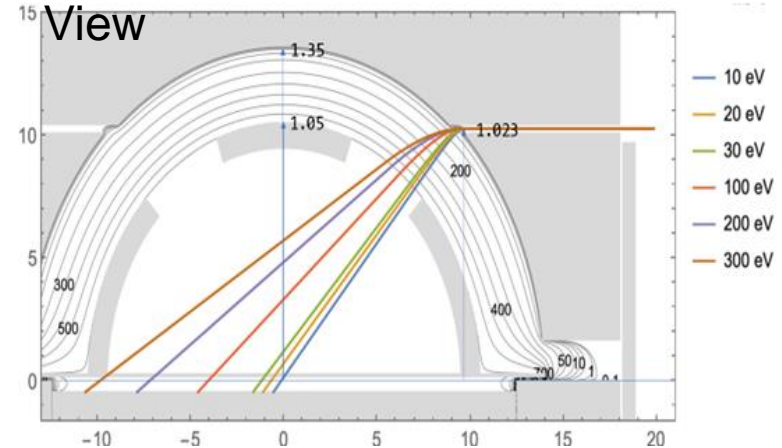
# Particle Analyzers



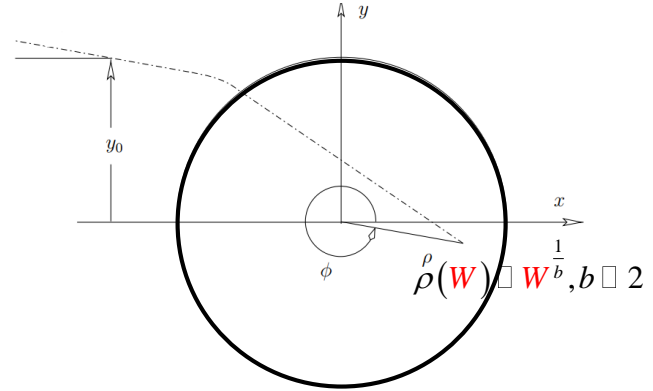
# Whalen Analyzer: Point-Spread Function



# Whalen Analyzer Side-View

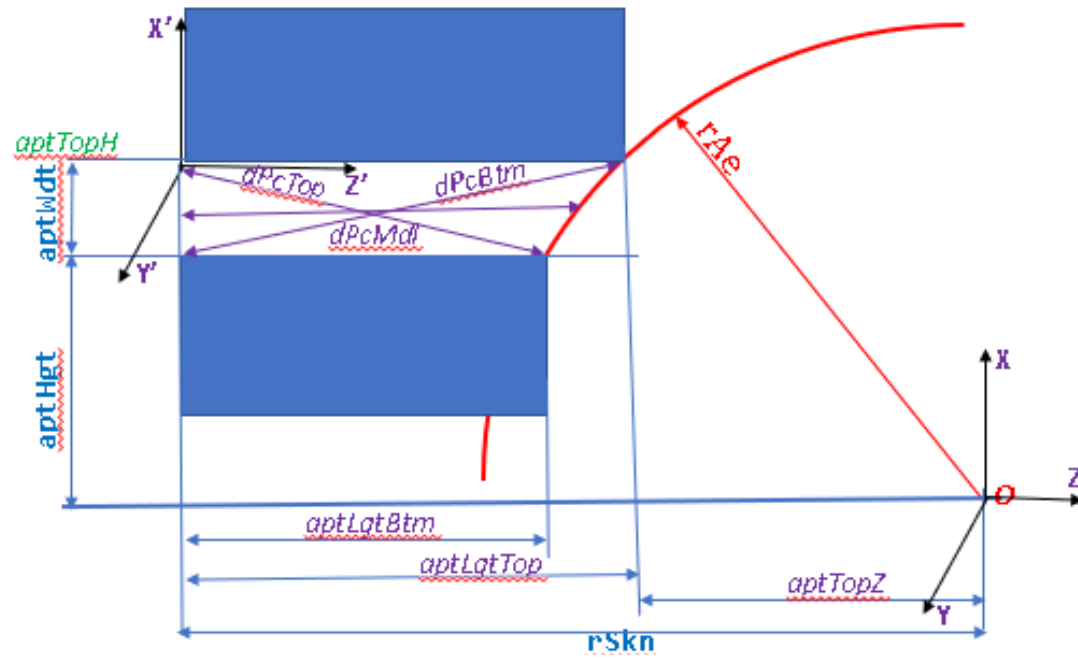


# Whalen Analyzer Top-View

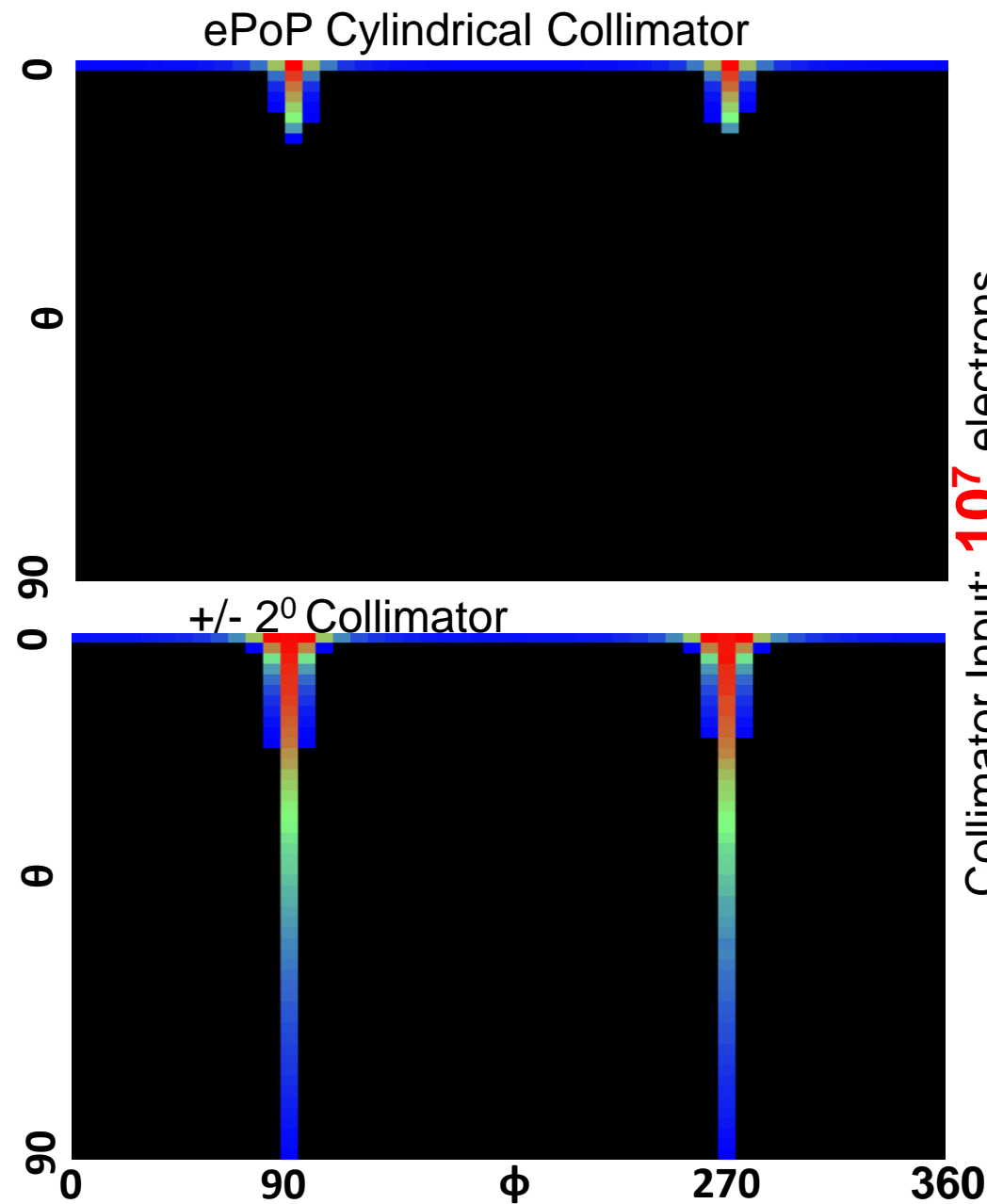
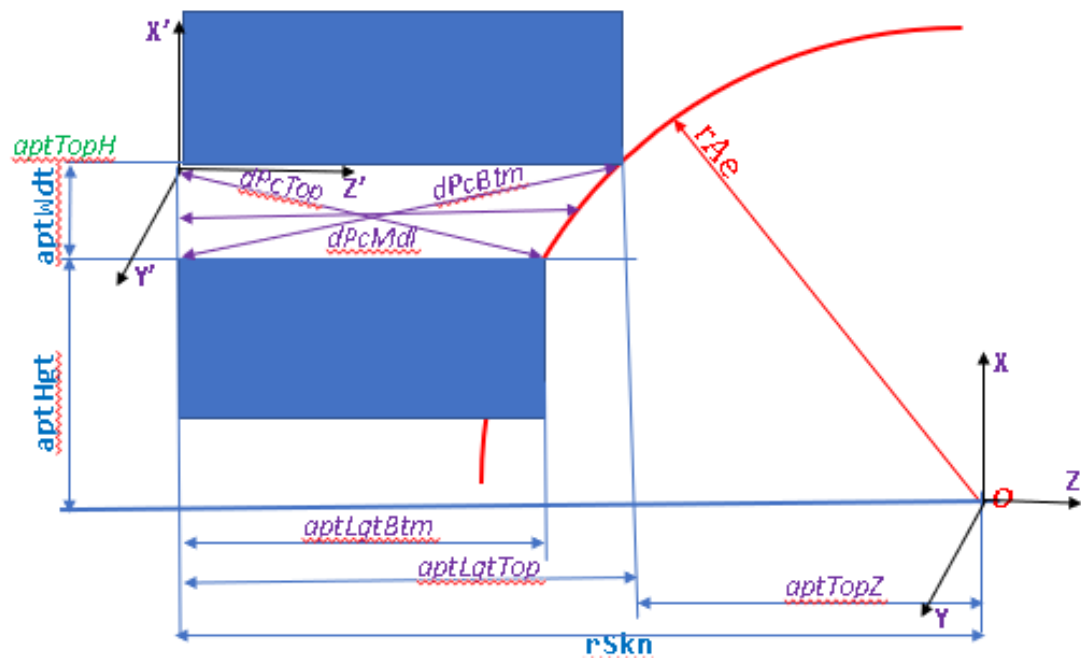


Side-view section (x - z plane) with on-axis particle trajectories of different energies.

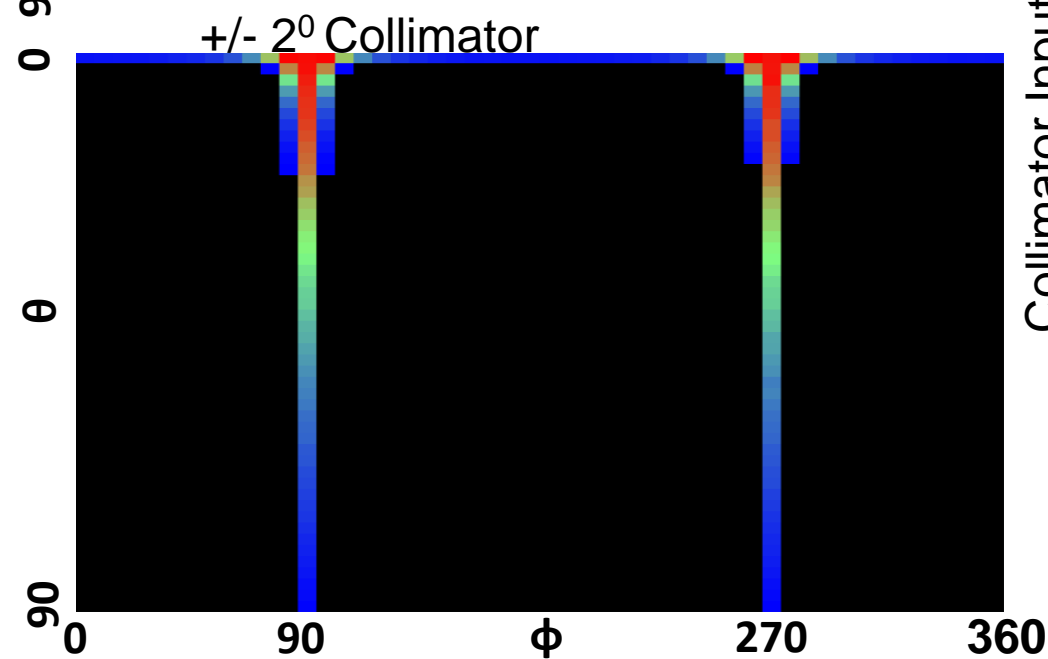
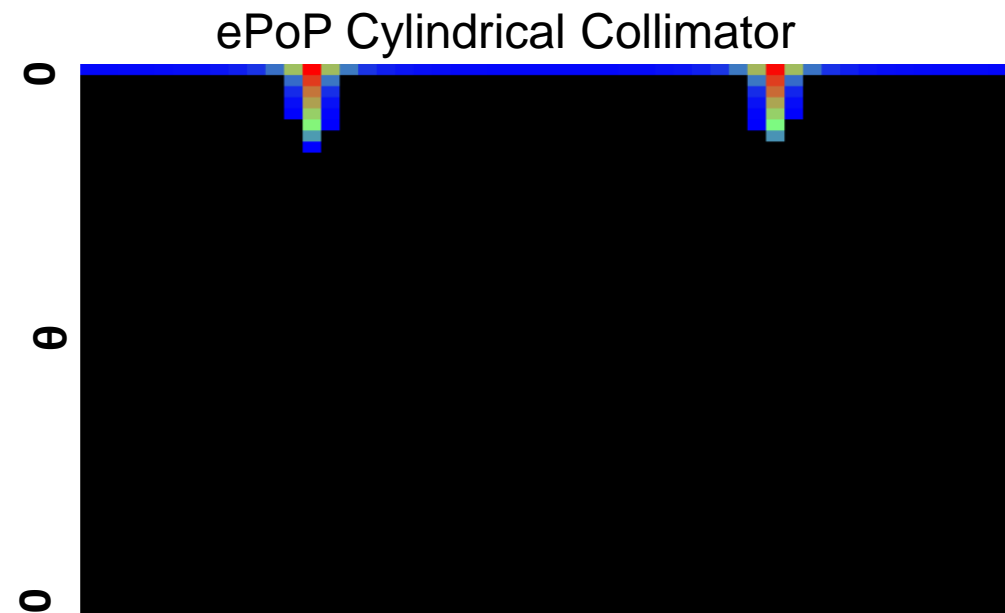
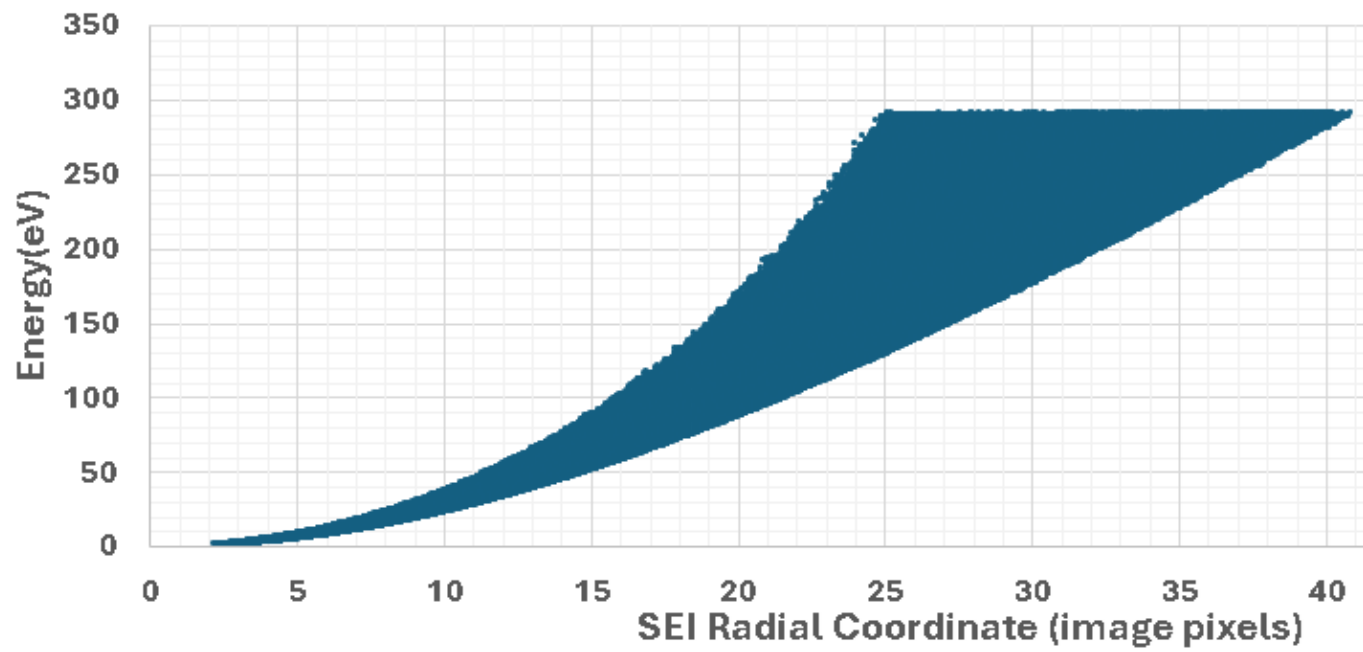
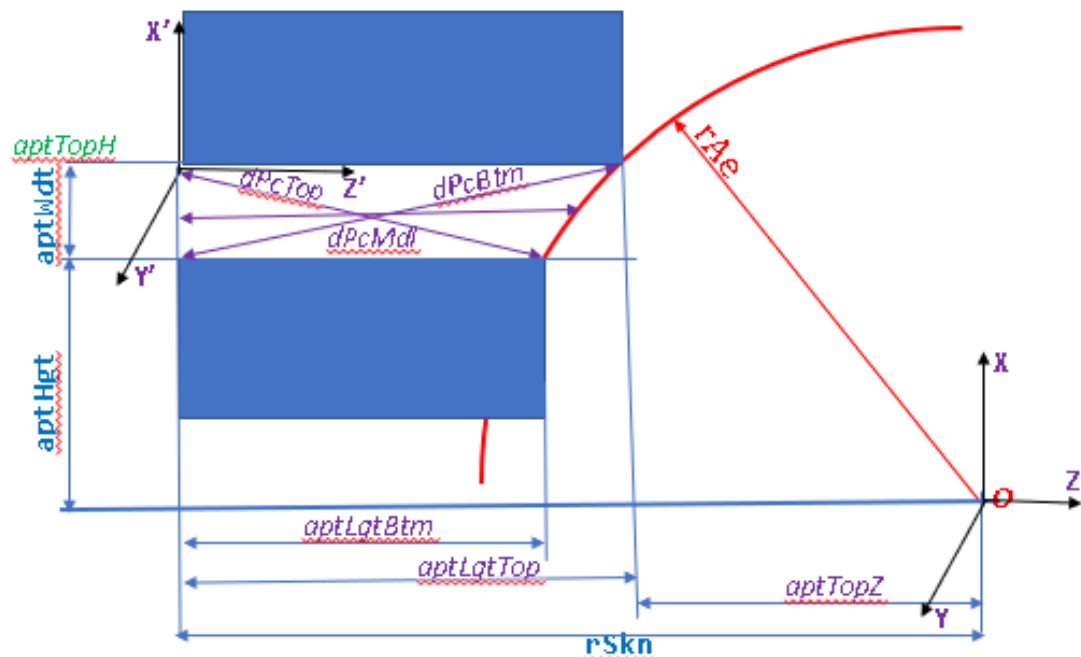
## Collimator Angular Distributions at the Detector Aperture



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Collimator Input: **10<sup>7</sup>** electronsOutput: **109,099** electronsOutput: **482,929** electrons

# SEI Shielding in the Strongly Coupled Plasmas

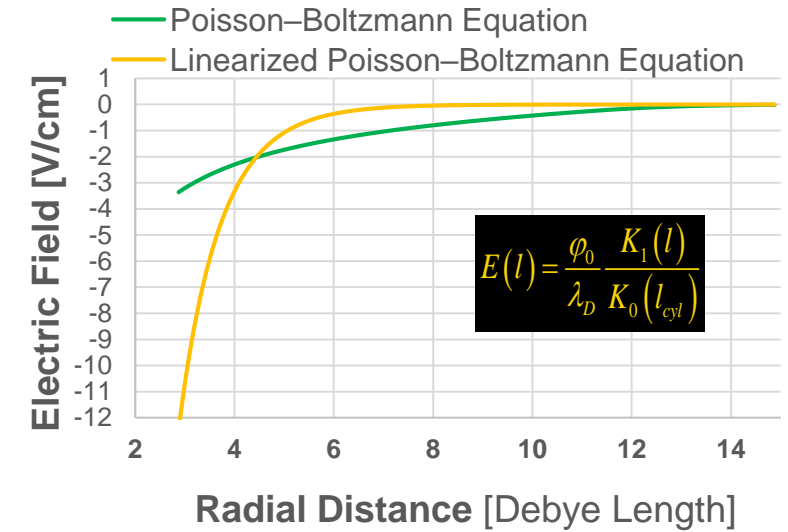
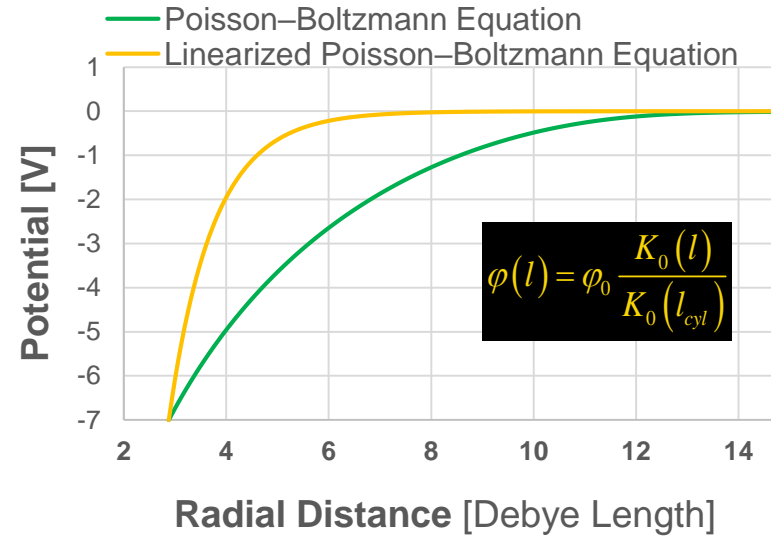
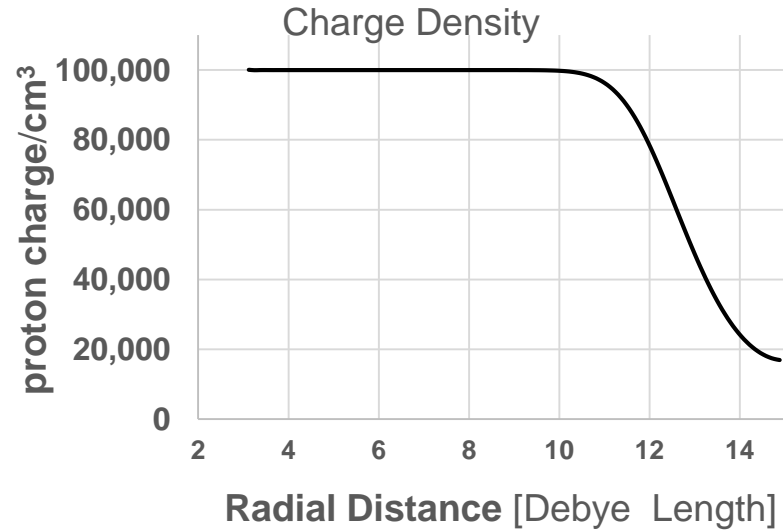
Poisson–Boltzmann equation in vector form:

$$\nabla^2 \varphi(\vec{r}) = 4\pi e n_0 (\exp(\chi) - 1), \chi = \frac{e\varphi(\vec{r})}{k_B T} < 0$$

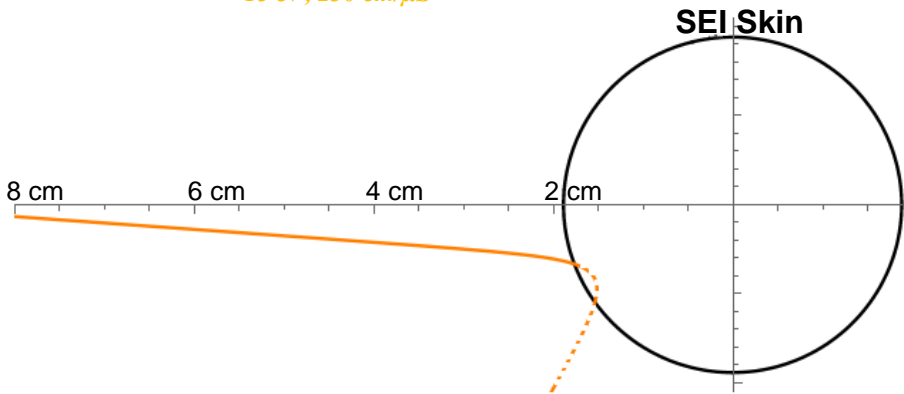
Parameter	Value
Electron Temperature	900 K
Plasma Density	100,000 cm <sup>-3</sup>
Debye Length	0.6547 cm
Cylinder (SEI) Radius	1.8825 cm
Cylinder Potential	-7 Volt

Poisson–Boltzmann equation in cylindrical coordinates:

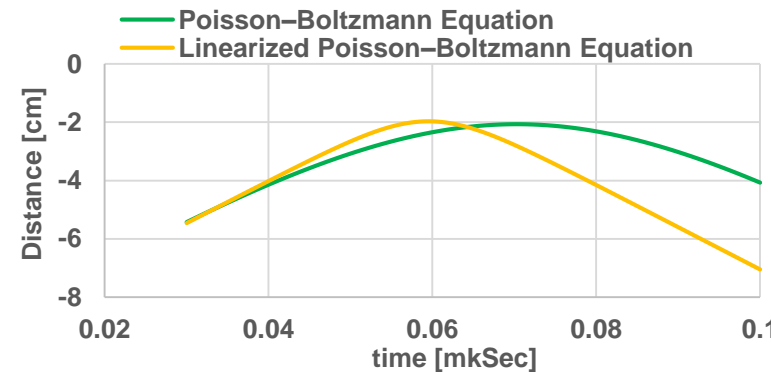
$$\frac{d^2 \chi(l)}{dl^2} + \frac{1}{l} \frac{d\chi(l)}{dl} = e^{\chi(l)} - 1; l = r / \sqrt{\frac{k_B T}{4\pi e^2 n_0}}$$



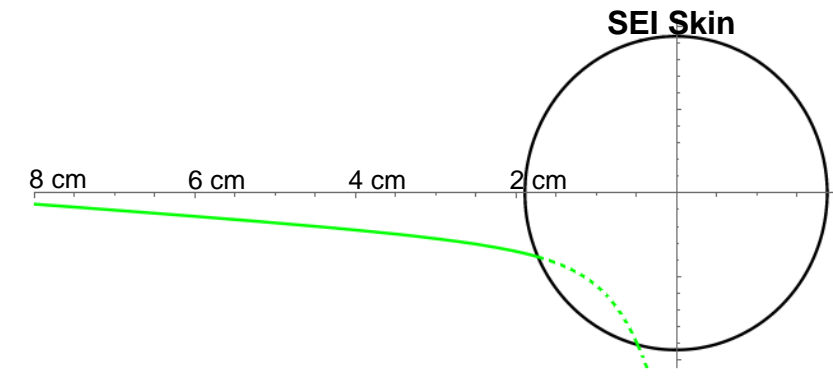
Linearized Poisson–Boltzmann Equation  
15 eV, 230 cm/μS



Electron Trajectories in the SEI Electric Field  
7 eV, 157 cm/μS



Exact Poisson–Boltzmann Equation  
15 eV, 230 cm/μS





# Electrons Deflection by the SEI Skin External Electric Field

**Conservation of Density in Phase Space, or the Liouville Theorem:** The distribution function is **constant** along any trajectory in phase space in the incompressible dynamical systems.

**Lemma:** Electron probability to hit the detector  $Eff_{cyl}$  for the **isotropic** incoming electron flux in the **infinite** cylindrical geometry depends on the electron initial **energy**  $w$  and the probe **potential**  $\varphi$  (in respect to the ambient plasma) and **does not depend** on the **shape** of the electric field:

$$Eff_{cyl} = \frac{\int I_{\hat{\Omega}}(w, \varphi) \vec{j}(\hat{\Omega}) \cdot \hat{n} d\hat{\Omega}}{\int_{2\pi} \vec{j}(\hat{\Omega}) \cdot \hat{n} d\hat{\Omega}} = \dots = 1 - \frac{e \cdot \varphi}{w}, w \in [e \cdot \varphi, \infty]$$

where  $\hat{\Omega}$  is a unit vector of an electron direction,  $\vec{j}$  is an external electron current density entering a cylindrical region,  $\hat{n}$  is a normal vector to the surface of a cylinder, and  $I_{\hat{\Omega}}(w, \varphi)$  is an indicator function that maps elements of directions to **one** if electron with energy  $w$  hit the detector in the presence of the electric field, and **zero** for the rest of directions.

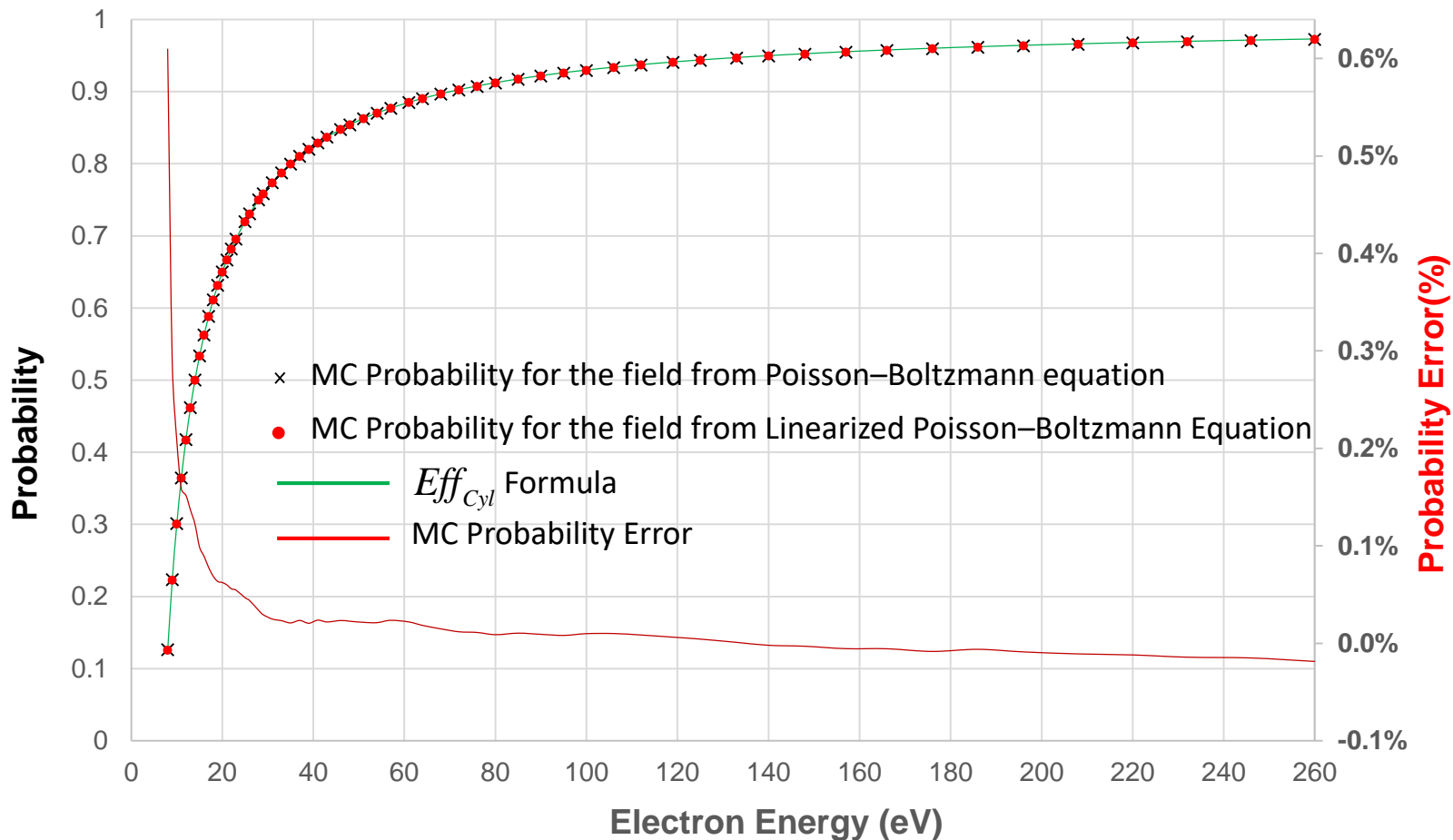
**Prove Idea:** The derivation for the probability to hit the detector assumes that any final direction at the collimator entrance at a given energy has the same probability (or distribution, according to the Liouville Theorem), except **prohibited** final directions, when electron can't penetrate sensor's electric field if it's traced back to space.

Compare with Mott-Smith and Langmuir, 1926, "The Theory of Collectors In Gaseous Discharge":

*"Special properties of the Maxwellian distribution: If the sheath has axial symmetry so that the equipotentials are coaxial circular cylinders, it is found from simple mechanical principles that the condition for a positive or a negative ion to reach the collector depends not upon the nature of the field of force along the whole orbit, but only upon the initial and final potentials and the initial velocity of the ion on entering the sheath".*

# MC Simulations of Electron Probability to Hit Detector

7



MC Simulation Parameters	Value
Electron Temperature	900 K
Plasma Density	100,000 cm <sup>-3</sup>
Debye Length	0.6547 cm
Cylinder (SEI) Radius	1.8825 cm
Cylinder Potential	-7 Volt

Leap-Frog trajectory calculations of isotopically distributed electrons coming from infinity and propagation in the radial electric field of the infinite right cylinder with the radius 1.8825 cm, electric potential -7[V], and plasma parameters from the table. Ten million electrons for each of 59 energies ranging from 8 to 260 eV were traced for two electric field models applied: One for the solution of the Poisson–Boltzmann equation, and one for the solution of the linearized Poisson–Boltzmann equation.

# Conclusions

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*\*\*\* Plasma Factory \*\*\**

*Done.*

*Not Covered in DASP-2024*

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*\*\*\* SEI Point-Spread Function \*\*\**

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*Will be Presented Tomorrow*

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*\*\*\* SEI Photoelectron Spectra Deconvolution \*\*\**

*Work in progress:*

*Satellite Potential Needed ...*

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## 4. A photoelectron model may help to explain the need in practice for a large negative faceplate bias to enable estimation of ion density

*\*\*\* Satellite Instrument Shielding in the Strongly Coupled Plasma \*\*\**

*Work in progress:*

*Satellite Potential Needed ...*